

## *Original Investigations*

# **Transferability of Some Properties of Localized Molecular Orbitals**

## **II. Electric Moments**

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The transferability of the first and second order electric moments of localized orbitals have been shown. Some moment characteristics have also been investigated. The standard deviation of first order localized moments is less than 2.0%. The standard deviation of second order electric moment components are larger: although they do not exceed 4.0% for bond orbitals, for lone pair ones they vary about 1.5–10.1%.

**Key words:** Localized MO's, Transferability of ~.

## **1. Introduction**

In a series of papers we have investigated the various properties of localized charge densities [1–4]. The localized orbitals calculated by using basis sets including polarization functions have also been analyzed [5]. Then the transferable property of some quantities connected with localized charge distributions have been studied (Part I, [6]). The following ten and eighteen-electron systems were included in the investigation:

CH<sub>4</sub>, CH<sub>3</sub>CH<sub>3</sub>, CH<sub>3</sub>NH<sub>2</sub>, CH<sub>3</sub>OH,

NH<sub>3</sub>, NH<sub>2</sub>NH<sub>2</sub>, NH<sub>2</sub>OH,

H<sub>2</sub>O, H<sub>2</sub>O<sub>2</sub>.

A uniform double-zeta basis set was used (more detail see in Refs. [1, 2]). The localized orbitals were obtained by the Edmiston–Ruedenberg procedure [7]. The actual geometries and the calculated total energies can be found in Part I, where we discussed the extent of transferability for the energy quantities connected with localized orbitals.

In the present paper the transferability of the electric moments of localized orbitals are investigated. The electric charge distribution in a molecule is a characteristic of the molecule itself [8, 9]. For many purposes, however, a detailed knowledge of the charge density is not necessary. Instead, the electric multipole moments up to certain order (third or fourth) will suffice. The charge distribution of a molecule can be described using the electric moments of its LMO density.

## 2. First Order Electric Moments

The first order electric moment components for a localized orbital  $\varphi_k$  is defined as follows:

$$r_k(u) = \langle \varphi_k | u | \varphi_k \rangle \quad u = x, y, z.$$

The center of the charge distribution is situated at the endpoint of vector  $\mathbf{r}_k$ , and  $|\mathbf{r}_k|$ , called centroid length, is:

$$|\mathbf{r}_k| = [\sum r_k^2(u)]^{1/2} \quad u = x, y, z.$$

The origin of vector  $\mathbf{r}_k$  is at the nucleus of the central atom. The non-vanishing total molecular dipole moments obtained for the 18-electron systems studied are the following in atomic unit:

CH<sub>3</sub>NH<sub>2</sub>: 0.707  
 CH<sub>3</sub>OH: 0.901  
 NH<sub>2</sub>OH: 0.392  
 NH<sub>2</sub>NH<sub>2</sub>: 1.036  
 H<sub>2</sub>O<sub>2</sub>: 1.006

The values are expected to be close to those which can be obtained in the (sp) limit.

The first order electric moments for bond and lone pair localized orbitals in the systems considered are given in Table 1. Several regularities could be found. First, it is remarkable that the standard deviation for the centroid lengths of bond pair orbitals is not larger than 1.0%, while that for lone pair orbitals is about 1.5–2.0%. This means that the bond and lone pair localized moments seem to be transferable to a rather good approximation. It should be noted – as also mentioned in Part I – that the presence of polarization functions is expected to cause only small systematic changes in the overall transferability of the moments. Similarly to the energy contributions (see Part I [6]), the deviation between the centroid lengths of orbitals of a given type in the same molecule is not larger, than the deviation of the centroid lengths for same type of orbitals in different molecules.

It is also remarkable that the standard deviation for centroid lengths is decreasing both for bond and lone pair orbitals with increasing nuclear charge of the central heavy atom. From the results one can see that not only the

**Table 1.** Centroid lengths (in a.u.)

	<b>Bond CH</b>	
CH <sub>4</sub>	1.36130	
CH <sub>3</sub> CH <sub>3</sub>	1.37756	
CH <sub>3</sub> NH <sub>2</sub> (1)	1.40337	
(2)	1.37265	
CH <sub>3</sub> OH (1)	1.39330	
(2)	1.37674	
Mean value	1.38052	
St. dev. (N%)	1.00	
	<b>Bond NH</b>	<b>Lone p.N</b>
NH <sub>3</sub>	1.14778	0.71288
NH <sub>2</sub> NH <sub>2</sub> (1)	1.16823	0.71155
(2)	1.14494	—
CH <sub>3</sub> NH <sub>2</sub>	1.15416	0.70115
NH <sub>2</sub> OH	1.15292	0.73843
Mean value	1.15360	0.71603
St. dev. (N%)	0.70	1.91
	<b>Bond OH</b>	<b>Lone p.O</b>
H <sub>2</sub> O	0.97796	0.60525
H <sub>2</sub> O <sub>2</sub> (1)	0.97130	0.62614
(2)	—	0.62933
CH <sub>3</sub> OH	0.97946	0.60306
NH <sub>2</sub> OH	0.97339	0.62153
Mean value	0.97555	0.61706
St. dev. (N%)	0.34	1.75

standard deviations are small, but the values of centroid length for the series of studied C, N and O containing compounds are well-separated as well. This property can be useful in characterizing localized orbitals by their electric moments even for larger systems.

### 3. Second Order Electric Moments

The second order electric moment components for localized orbital  $\varphi_k$  can be defined in the following way:

$$Q_k(u, v) = \langle \varphi_k | uv | \varphi_k \rangle \quad u, v = x, y, z.$$

The origin of second moment tensor is usually shifted into the endpoint of the corresponding centroid vector. In this case the second moment tensor has the components

$$\bar{Q}_k(u, v) = \langle \varphi_k | [u - r_k(u)][v - r_k(v)] | \varphi_k \rangle \quad u, v = x, y, z.$$

In this section we analyze the shifted (located) second order electric moments for the systems considered.

The total quadrupole moments of the compounds studied are presented in Table 2. The values tabulated are given in atomic units. The eigenvalues –

**Table 2.** Principal values of the diagonalized quadrupole moment tensors (calculated relative to the center of mass)

	XX	YY	ZZ
CH <sub>3</sub> CH <sub>3</sub>	0.524	0.524	-1.048
CH <sub>3</sub> NH <sub>2</sub>	7.538	-8.079	0.541
CH <sub>3</sub> OH	4.510	-2.873	-1.637
NH <sub>2</sub> NH <sub>2</sub>	3.626	-3.726	0.100
NH <sub>2</sub> OH	5.502	-4.794	-0.708
H <sub>2</sub> O <sub>2</sub>	4.380	-3.865	-0.515

which characterize the second moment tensor for the individual localized orbital  $\varphi_k$  – can be obtained by diagonalizing  $\bar{\mathbf{Q}}_k$ . The diagonalized tensor has the components  $q_x^2(k)$ ,  $q_y^2(k)$ ,  $q_z^2(k)$ . The quantities  $q_x$ ,  $q_y$  and  $q_z$  are called the dispersions in the given direction (for detail see [2]). If  $\bar{\mathbf{Q}}_k$  has cylindrical symmetry the  $q_y^2 = q_z^2$  equation will hold.

The second order electric moments for the bond orbitals are given in Table 3.a. As to the form of tensors  $\mathbf{Q}$  of bond pair charge densities it can be seen that  $q_x \gg q_y > q_z$  in all cases. When the symmetry allows (CH<sub>4</sub>, e.g.) then  $q_y = q_z$ , in all other cases  $q_y$  and  $q_z$  are different but their deviation does not exceed 5%

**Table 3.a.** Second order electric moments for bond orbitals<sup>a</sup>

Bond CH	$q_x^2$	$q_y^2$	$q_z^2$
CH <sub>4</sub>	1.25470	0.70290	0.70290
CH <sub>3</sub> CH <sub>3</sub>	1.20412	0.69961	0.69461
CH <sub>3</sub> NH <sub>2</sub> (1)	1.14860	0.69156	0.68365
(2)	1.17881	0.68657	0.68233
CH <sub>3</sub> OH (1)	1.12994	0.69037	0.67353
(2)	1.14624	0.68592	0.66291
Mean value	1.17705	0.69268	0.68330
St. dev. (N%)	3.59	0.85	1.91
Bond NH	$q_x^2$	$q_y^2$	$q_z^2$
NH <sub>3</sub>	1.05596	0.55113	0.54942
NH <sub>2</sub> NH <sub>2</sub> (1)	0.99822	0.55259	0.54725
(2)	0.99634	0.55030	0.54060
CH <sub>3</sub> NH <sub>2</sub>	1.03901	0.56165	0.54635
NH <sub>2</sub> OH	0.97669	0.54240	0.53169
Mean value	1.01324	0.55160	0.54306
St. dev. (N%)	2.91	1.11	1.18
Bond OH	$q_x^2$	$q_y^2$	$q_z^2$
H <sub>2</sub> O	0.92811	0.44453	0.44256
H <sub>2</sub> O <sub>2</sub>	0.85446	0.43046	0.42671
CH <sub>3</sub> OH	0.90814	0.46797	0.44392
NH <sub>2</sub> OH	0.88048	0.43734	0.43637
Mean value	0.89280	0.44508	0.43740
St. dev. (N%)	3.12	3.18	1:55

<sup>a</sup> in (a.u.)<sup>2</sup>.

**Table 3.b.** Second order electric moments for lone pair orbitals<sup>a</sup>

Lone p.N	$q_x^2$	$q_y^2$	$q_z^2$
NH <sub>3</sub>	0.90192	0.60622	0.60622
NH <sub>2</sub> NH <sub>2</sub>	0.86289	0.62319	0.57990
CH <sub>3</sub> NH <sub>2</sub>	0.93620	0.64952	0.59659
NH <sub>2</sub> OH	0.77492	0.59151	0.58604
Mean value	0.86898	0.61760	0.59218
St. dev. (N%)	6.93	3.49	1.70
Lone p.O	$q_x^2$	$q_y^2$	$q_z^2$
H <sub>2</sub> O	0.70506	0.45887	0.45628
H <sub>2</sub> O <sub>2</sub> (1)	0.60074	0.45187	0.43914
(2)	0.60248	0.45663	0.44479
CH <sub>3</sub> OH	0.69506	0.52082	0.45871
NH <sub>2</sub> OH	0.63293	0.48035	0.44826
Mean value	0.64726	0.47370	0.44944
St. dev. (N%)	10.10	5.38	1.61

<sup>a</sup> in (a.u.)<sup>2</sup>.

(see Table 3.a). The standard deviation of components  $q_x^2$  for bonds CH, NH and OH vary between 2.9 and 4.6%. i.e., the deviations are much larger than those found for the first order localized moments. The standard deviations for the components  $q_y^2$  and  $q_z^2$  are smaller than those for  $q_x^2$  in all of the three series.

It is interesting, that for the CH bond the standard deviation of  $q_y^2$  is smaller than that of  $q_z^2$ , while for the NH bond they are nearly equal, and for the OH orbitals the standard deviation of components  $q_y^2$  is much larger than that of  $q_z^2$ .

The second order electric moments for lone pair orbitals differ more within the same series from each other than those for bond orbitals. The results are presented in Table 3.b. The standard deviation for components  $q_x^2$  for O containing series exceeds even 10.00% and the  $q_y^2$  and  $q_z^2$  are also rather large, varying between 1.6 and 5.4%. It is remarkable that while the components  $q_x^2$  are always larger for bonds than for lone pair orbitals in the system studied, it does not hold the same for the values of  $q_y^2$  and  $q_z^2$  (see, e.g., NH bond and OH bond in Table 3.a, or N lone pair and O lone pair orbitals in Table 3.b).

The regularity generally holds that the first order electric moments for any type of localized orbitals are less sensitive to the change of molecular environment than the second order ones.

#### 4. Moment Characteristics

Some quantities have been proposed recently [2] in order to characterize the charge distribution of bond and lone pair orbitals. First, the angles between the centroid vectors of bond orbitals are investigated. They are in all cases smaller than the corresponding valence angles. The deviation of bond centroid vector from the valence axis in degrees are given (under the heading  $\alpha$ ) in Table 4. The  $\alpha$  values do not exceed the 3.0° for any of bond orbitals.

**Table 4.** Deviation angles for the centroid vectors (in degrees). Definition of  $\alpha$  and  $\beta$  see in text

	<b>Bond</b>	<b>CH</b>		
	$\alpha$	$\beta$		
CH <sub>4</sub>	0.0	0.0		
CH <sub>3</sub> CH <sub>3</sub>	0.4	0.8		
CH <sub>3</sub> NH <sub>2</sub> (1)	0.2	2.1		
(2)	0.5	1.0		
CH <sub>3</sub> OH (1)	0.7	1.6		
(2)	0.4	3.4		
Mean value	0.4	1.5		
St. dev. (N%)	60.3	72.7		
	<b>Bond</b>	<b>NH</b>	<b>Lone</b>	<b>p.N</b>
	$\alpha$	$\beta$	$\alpha$	$\beta$
NH <sub>3</sub>	0.6	3.2	—	8.5
NH <sub>2</sub> NH <sub>2</sub> (1)	2.2	4.5	73.2	4.1
(2)	2.7	5.1	—	—
CH <sub>3</sub> NH <sub>2</sub>	1.7	5.2	73.2	11.0
NH <sub>2</sub> OH	1.4	5.0	70.9	10.1
Mean value	1.7	4.6	72.4	8.4
St. dev. (N%)	41.5	16.1	1.5	31.5
	<b>Bond</b>	<b>OH</b>	<b>Lone</b>	<b>p.O</b>
	$\alpha$	$\beta$	$\alpha$	$\beta$
H <sub>2</sub> O	0.6	2.4	—	6.4
H <sub>2</sub> O <sub>2</sub> (1)	0.6	1.9	75.4	10.8
(2)	—	—	76.3	11.7
CH <sub>3</sub> OH	1.6	5.2	77.1	11.7
NH <sub>2</sub> OH	0.5	7.2	74.1	10.1
Mean value	0.8	4.2	75.7	10.1
St. dev. (N%)	54.5	51.6	1.5	19.4

In cases of lone pair orbitals under the heading  $\alpha$  the deviation in degree of the given centroid vector and the molecular axis (i.e., the axis which connects the two heavy atoms) are presented. The standard deviation for both the nitrogen and oxygen lone pair orbitals is about 1.5%. The deviation angles are only slightly larger for the oxygen than for the nitrogen lone pair orbitals. The small standard deviations suggest that the orientation of lone pair centroid vectors are well defined in the 18-electron systems. Apart from the symmetric cases, the direction of major dispersion axis ( $q_x$ ), which has its origin at the center of charge of the given localized orbital, does not coincide with the corresponding centroid vector. These non-zero angles are presented in Table 4 under the heading  $\beta$ . The standard deviation for the angles  $\beta$  are rather large both for bond and lone pair orbitals. The deviation angles  $\beta$  for a system containing similar orbitals were expected to be close to each other (see, e.g., NH<sub>2</sub>NH<sub>2</sub>(1) and (2) or H<sub>2</sub>O<sub>2</sub> (1) and (2)).

In addition to the deviation angles, some further quantities have been defined for characterizing the main differences between bond and lone pair charge

distributions (for all definitions see [2]). The asymmetry of bond orbitals ( $A_{\text{orb}}$ ) is evidently decreasing as the nuclear charge of central heavy atom increases. The values are given in Table 5. The results parallel with those found for the centroid lengths—i.e. the CH, NH or OH bonds could be identified even by their LMO symmetry. It does not hold the same for the lone pair orbitals, as they vary between 0.72 and 0.84 for both nitrogen and oxygen containing systems.

The further moment characteristics, namely the effective solid angles ( $\Omega_{\text{eff}}$ ), the dispersion ratios ( $R_d$ ) and dispersion products ( $P_d$ ) of localized charge densities show similar regularities than the orbital asymmetries do. All of the three quantities form well-separated groups for CH, NH and OH bond orbitals (see Table 5). The values of two characteristics for the lone pair orbitals, however, do overlap in the nitrogen and oxygen containing series: the  $\Omega_{\text{eff}}$  angles vary

**Table 5.** Moment characteristics

	$A_{\text{orb}}$	$\Omega_{\text{eff}}$ (sterad)	$R_d$	$P_d$ (a.u.) <sup>3</sup>
<b>Bond CH</b>				
CH <sub>4</sub>	1.215	0.933	1.336	0.787
CH <sub>3</sub> CH <sub>3</sub>	1.256	0.908	1.316	0.763
CH <sub>3</sub> NH <sub>2</sub> (1)	1.310	0.873	1.293	0.737
(2)	1.264	0.902	1.312	0.743
CH <sub>3</sub> OH (1)	1.311	0.878	1.287	0.725
(2)	1.286	0.887	1.304	0.722
<b>Bond NH</b>				
NH <sub>3</sub>	1.117	1.006	1.385	0.566
NH <sub>2</sub> NH <sub>2</sub> (1)	1.169	0.978	1.347	0.549
(2)	1.147	1.003	1.352	0.544
CH <sub>3</sub> NH <sub>2</sub>	1.132	1.003	1.370	0.565
NH <sub>2</sub> OH	1.166	0.981	1.349	0.531
<b>Bond OH</b>				
H <sub>2</sub> O	1.015	1.090	1.447	0.427
H <sub>2</sub> O <sub>2</sub>	1.051	1.073	1.412	0.396
CH <sub>3</sub> OH	1.028	1.110	1.411	0.434
NH <sub>2</sub> OH	1.037	1.085	1.420	0.410
<b>Lone p.N</b>				
NH <sub>3</sub>	0.751	2.040	1.220	0.576
NH <sub>2</sub> NH <sub>2</sub>	0.766	2.035	1.198	0.558
CH <sub>3</sub> NH <sub>2</sub>	0.725	2.108	1.226	0.602
NH <sub>2</sub> OH	0.839	1.925	1.147	0.518
<b>Lone p.O</b>				
H <sub>2</sub> O	0.721	2.103	1.237	0.384
H <sub>2</sub> O <sub>2</sub> (1)	0.808	1.985	1.176	0.337
(2)	0.810	1.988	1.156	0.350
CH <sub>3</sub> OH	0.723	2.180	1.192	0.408
NH <sub>2</sub> OH	0.781	2.049	1.168	0.369

about 2.0 degree while the  $R_d$  ratios between 1.14 and 1.24. These results suggest that the charge distribution of lone pair orbitals have similar "form" for all systems considered, i.e. they are not too much dependent on the central atom nuclear charge or on the molecular environment. The dispersion product, however, even for lone pair orbitals is rather different in the series of oxygen and nitrogen containing molecules. The values are about 5.1–6.0 for the former, while 3.3–4.1 for the latter ones. This result confirms the regularity found earlier: while the "form" of lone pair orbital charge densities are similar, the "region" which they occupy are the larger the smaller is the nuclear charge of the central heavy atom.

The conclusion can be drawn that the moment characteristics for lone pair orbitals as well as for bonds are transferable to a rather good approximation.

### 5. Charge Displacements in the LMO's of 18-Electron Systems

The LMO's of the systems considered are slightly delocalized extending into regions away from the center of their charge clouds. We have found that in the LMO's of the valence shell beyond the contributions from the principal components (two-center bonds or one-center lone pairs) those from the basis functions of vicinal H atoms are also appreciable. Similar results were obtained by England *et al.* for acyclic hydrocarbons [10]. The comparison of our results with theirs is, however, not straightforward because they use INDO approximation with a minimal hybrid basis.

Our results have also shown that in the LMO of a given bond or a lone pair the coefficients of basis functions of a vicinal H atom are the largest if the H atom is coplanar (or nearly coplanar) to the principal components. In these cases the magnitude of some of the coefficients usually exceeds the value of 0.05. If the vicinal H atom is in the cis position the coefficients of its basis functions and those of the principal components are of opposite sign. If the vicinal H atom is in trans position, the coefficients of its basis functions in the corresponding LMO are of the same sign as those of the principal components. For vicinal H atoms in gauche position the magnitudes of the coefficients are usually smaller than 0.05. The position of the vicinal H atom manifests itself in the magnitude of the centroid length of the corresponding LMO. For example, in  $\text{CH}_3\text{NH}_2$  the two equivalent CH bonds (denoted by (2) in Table 1) have nearly coplanar vicinal H atoms in trans position. As a consequence their centroid lengths are smaller than that of the other CH bond (denoted by (1) in Table 1). Similar situation occurs in  $\text{NH}_2\text{NH}_2$  where the centroid length of the two NH bonds heaving nearly coplanar vicinal H atoms (denoted by (2) in Table 1) is smaller than that denoted by (1). In  $\text{CH}_3\text{OH}$  the centroid length of those CH bond is smaller (denoted by (2) in Table 1) which have a coplanar vicinal H atom in cis position.

The results obtained on the study of the charge displacement in the LMO's of 18-electron systems provided further information on the spatial distribution of localized charge densities (see, e.g., in [4]) in various molecules.



## 6. Conclusion

Several regularities have been found as investigating the localized moments and moment characteristics for some related ten- and eighteen-electron systems. The first order electric moments are transferable to a rather good approximation. As to the values of second order electric moments, they are more sensitive to the change of molecular environment. It is expected that results obtained for these quantities by the use of larger basis sets (including polarization functions) would not alter very much these conclusions.

The study of moment characteristics has shown that, in addition to their advantages pointed out in [2], they are even transferable to a fairly good approximation.

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